

Analysis of a hybrid p-Multigrid method for the discontinuous Galerkin discretisation of the Euler equations

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Eccomas CFD - 5 september 2006

Research supported by the Walloon Region and the European funds ERDF and ESF



Acknowledgments

Joint work CENAERO - UCL (GCE)

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- Dr. N. Cheveaugeon(UCL - GCE)
- Ir. P.-E. Bernard (UCL - MEMA)
- Dr. P. Geuzaine (CENAERO)

Outline

- Motivation
- DGFEM
- Solvers
 - Newton-Krylov-ILU(0)
 - h- or p-Multigrid
 - ♦ General FAS framework
 - ♦ Theoretical analysis
- Comparison: flow around NACA0012
- Conclusions And Future Work

Motivation

Advantages of higher order DGFEM

- increased order improves ratio precision wrt RAM and CPU
- maintains first order stencil on “unstructured level”
 - parallelisability
 - low overhead associated to mesh
- data per element is locally “structured”
 - no global renumbering of unknowns
 - casting of global operator into matrix-matrix operations
- hp adaptivity

Shortfalls for the moment

- efficient solution methods ?
- expensive evaluation of the terms due to Gauss quadrature
- viscous terms
- turbulence model
- shock capturing

DGFEM

Find the solution u in \mathcal{D} to

$$\mathcal{L}(u) = \nabla \cdot \vec{f}(u) = 0$$

with appropriate boundary conditions.

Approximate solution in broken space \mathcal{U}

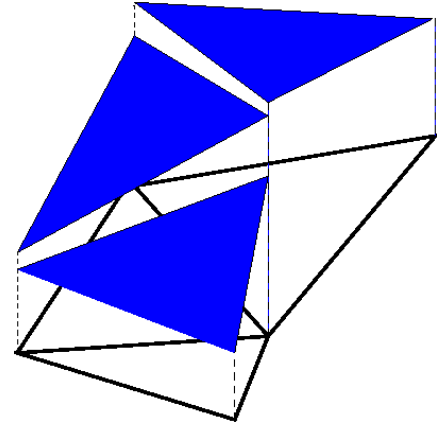
$$\tilde{u} = \sum_i \mathbf{u}_i \phi_i, \phi_i \in \mathcal{U}$$

Galerkin weighting:

$$\mathbf{r}_j = \int_{\mathcal{D}} \phi_j \frac{\partial \tilde{u}}{\partial t} dV + \int_{\mathcal{D}} \phi_j \nabla \cdot \vec{f}(\tilde{u}) dV = 0, \forall \phi_j \in \mathcal{U}$$

$$\mathbf{r}_j = \int_T \phi_j \frac{\partial \tilde{u}}{\partial t} dV + \oint_{\partial T} \phi_j \vec{f}(\tilde{u}) \cdot \vec{n} dS - \int_T \nabla \phi_j \cdot \vec{f}(\tilde{u}) dV$$

$$= \int_T \phi_j \frac{\partial \tilde{u}}{\partial t} dV + \oint_{\partial T} \phi_j \mathcal{H}(\tilde{u}^+, \tilde{u}^-, \vec{n}) dS - \int_T \nabla \phi_j \cdot \vec{f}(\tilde{u}) dV$$



Solvers

- Runge-Kutta local timestepping

$$\Delta\tau = CFL \cdot \frac{L}{u} \cdot \frac{1}{2p+1}$$

- Newton-Krylov with ILU preconditioner
- Multigrid
 - mesh resolution : h-Multigrid
 - order of solution : p-Multigrid

Newton-Krylov-ILU

- Nonlinear iteration : Newton solution for damped equations

$$\mathbf{r}_j^* = \left(\frac{\tilde{\mathbf{u}}^n - \tilde{\mathbf{u}}^{n-1}}{\Delta\tau^n} + \nabla \cdot \vec{f}(\tilde{\mathbf{u}}^n), \phi_i \right) = 0, \forall \phi_i \in \mathcal{U}$$

resulting in

$$\left(\frac{\mathbf{M}}{\Delta\tau^n} + \frac{\partial \mathbf{r}}{\partial \mathbf{u}} \right) \cdot \Delta \mathbf{u}^n = -\mathbf{r}(\mathbf{u}^{n-1})$$

- Linear iterations : GMRES
 - “Matrix-free” implementation

$$\left(\frac{\mathbf{M}}{\Delta\tau^n} + \frac{\partial \mathbf{r}}{\partial \mathbf{u}} \right) \cdot \mathbf{s} \approx \frac{\mathbf{r}^*(\mathbf{u}^n + \epsilon \cdot \mathbf{s}) - \mathbf{r}^*(\mathbf{u}^n)}{\epsilon}$$

- block ILU preconditioner : one block per element, two per face
 - ♦ 2D : dofs/element $\sim p^2$ hence elementary matrix $\sim p^4$
 - ♦ 3D : dofs/element $\sim p^3$ hence elementary matrix $\sim p^6$

Multigrid

- MG uses successively coarser representations of the solution to iterate on
 - grid size: h-Multigrid (Bastian et al. 2000, Van der Vegt et al. 2002)
 - interpolation order: p-Multigrid (Helenbrook et al. 2003, Fidkowski et al. 2004)
- MG has optimal complexity (TME): ie. work needed to solve the system is only proportional to the number of dofs.
- MG convergence hinges upon
 - fast reduction of high-frequency errors - simple smoothers suffice
 - good quality of transfer operators
- MG is extremely easy to implement in a discontinuous interpolation framework

Multigrid : FAS for Galerkin FEM

“Fine” solution $\tilde{u}^p \in \mathcal{U}^p$

“Coarse” solution $\tilde{u}^q \in \mathcal{U}^q$

- presmoothing on level p , leading to solution \tilde{u}^{p*}
- restriction

$$\tilde{u}^{q*} = \mathcal{T}^{qp} \left(\tilde{u}^{p*} \right)$$

- defect correction

$$\left(\mathcal{L} \left(\tilde{u}^q \right) - \mathcal{L} \left(\tilde{u}^{q*} \right) + \mathcal{L} \left(\tilde{u}^{p*} \right), v^q \right) = 0 \quad \forall v^q \in \mathcal{U}^q$$

- prolongation

$$\tilde{u}^p = \tilde{u}^{p*} + \mathcal{T}^{pq} \left(\tilde{u}^q - \tilde{u}^{q*} \right)$$

- post-smoothing: perform additional iterations on level p to smooth the corrected solution

Multigrid : Solution transfer

\mathcal{T}^{ba} is defined as the L_2 projection from $u^a \in \mathcal{U}^a$ to $u^b \in \mathcal{U}^b$:

$$\sum_i (\phi_k^b, \phi_j^b) u_j^b = \sum_j (\phi_k^b, \phi_i^a) u_i^a, \quad \forall \phi_k^b \in \mathcal{U}^b$$

Equivalent discrete operator

$$\mathbf{u}^b = \mathbf{T}^{ba} \cdot \mathbf{u}^a$$

$$\mathbf{T}^{ba} = (\mathbf{M}^{bb})^{-1} \cdot \mathbf{M}^{ba}$$

where

$$\mathbf{M}_{ij}^{bb} = (\phi_i^b, \phi_j^b)$$

$$\mathbf{M}_{ij}^{ba} = (\phi_i^b, \phi_j^a)$$

Multigrid : Residual restriction

Avoid modified residual operators associated to mixed weighting

- project residual function on \mathcal{U}^p using L_2 projection:

$$\mathcal{L}(\tilde{u}^p) \approx \sum_i \mathbf{l}_i^p \phi_i^p$$

$$\sum_i \mathbf{l}_i^p (\phi_i^p, \phi_j^p) \approx (\mathcal{L}(\tilde{u}^p), \phi_j^p) = \mathbf{r}_j^p$$

- then Galerkin weighting becomes trivial:

$$(\mathcal{L}(\tilde{u}^p), \phi_j^q) \approx \sum_i \mathbf{l}_i^p (\phi_i^p, \phi_j^q)$$

Equivalent discrete operator:

$$\mathbf{r}^q = \tilde{\mathbf{T}}^{qp} \cdot \mathbf{r}^p$$

$$\tilde{\mathbf{T}}^{qp} = \mathbf{M}^{qp} \cdot (\mathbf{M}^p)^{-1} = (\mathbf{T}^{pq})^T$$

Multigrid : convergence analysis

Suppose $\mathcal{L}(\cdot)$ is a linear operator, then

$$\left(\mathcal{L}(\tilde{u}^p), \phi_i^p \right) = \left(f, \phi_i^p \right), \quad \forall \phi_i^p \in \mathcal{U}^p$$

is equivalent to

$$\mathbf{L}^p \cdot \mathbf{u}^p = \mathbf{f}^p$$

$$\mathbf{L}_{ij}^p = \left(\phi_i^p, \mathcal{L}(\phi_j^p) \right)$$

Suppose that interpolation space \mathcal{U}^q is “embedded” in \mathcal{U}^p

$$\phi_i^q = \alpha_{ij}^{qp} \cdot \phi_j^p, \quad \forall \phi_i^q \in \mathcal{U}^q$$

then

$$\alpha^{qp} = \left(\mathbf{M}^{qp} \cdot \left(\mathbf{M}^{pp} \right)^{-1} \right) = \tilde{\mathbf{T}}^{qp} = \left(\mathbf{T}^{pq} \right)^T$$

Multigrid : convergence analysis

We can prove in this case

$$\begin{aligned}\mathbf{L}_{ij}^q &= \left(\phi_i^q, \mathcal{L} \left(\phi_j^q \right) \right) \\ &= \alpha_{ik}^{qp} \cdot \left(\phi_k^p, \mathcal{L} \left(\phi_l^p \right) \right) \cdot \alpha_{jl}^{qp}\end{aligned}$$

Hence for the chosen transfer operators the discrete coarse grid approximation (**DCGA**) coincides with the Galerkin coarse grid approximation (**GCGA**) :

$$\mathbf{L}^q = \tilde{\mathbf{T}}^{qp} \cdot \mathbf{L}^p \cdot \mathbf{T}^{pq}$$

Hence the discretised defect correction equation

$$\mathbf{L}^q \cdot \left(\mathbf{u}^q - \mathbf{u}^{q*} \right) + \tilde{\mathbf{T}}^{qp} \cdot \left(\mathbf{L}^p \cdot \mathbf{u}^{p*} - \mathbf{f}^p \right) = 0$$

translates to:

$$\tilde{\mathbf{T}}^{qp} \left(\mathbf{L}^p \cdot \left(\mathbf{u}^{p*} + \mathbf{T}^{pq} \cdot \left(\mathbf{u}^q - \mathbf{u}^{q*} \right) \right) - \mathbf{f}^p \right) = 0$$

Multigrid : convergence analysis

We define the error e^p as:

$$\mathbf{L}^p \mathbf{e}^p = \mathbf{L}^p \mathbf{u}^p - \mathbf{f}^p = \mathbf{r}^p$$

The error e^p after one two-level iteration now becomes

$$\mathbf{e}^p = \left(\mathbf{I}_p - \mathbf{T}^{pq} \cdot \left(\mathbf{L}^q \right)^{-1} \cdot \tilde{\mathbf{T}}^{qp} \cdot \mathbf{L}^p \right) \cdot \mathbf{e}^{p*}$$

We may now decompose the error e^{p*}

$$\mathbf{e}^{p*} = \mathbf{e}_s^{p*} + \mathbf{e}_r^{p*}$$

$$\mathbf{e}_s^{p*} = \left(\mathbf{T}^{pq} \cdot \left(\tilde{\mathbf{T}}^{qp} \mathbf{T}^{pq} \right)^{-1} \cdot \tilde{\mathbf{T}}^{qp} \right) \cdot \mathbf{e}^{p*}$$

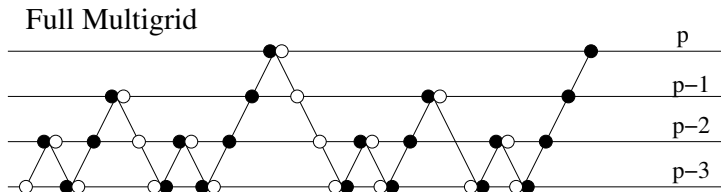
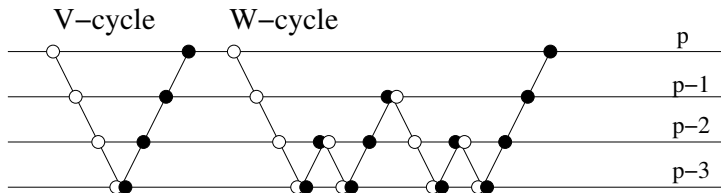
$$\mathbf{e}_r^{p*} = \left(\mathbf{I}_p - \mathbf{T}^{pq} \cdot \left(\tilde{\mathbf{T}}^{qp} \mathbf{T}^{pq} \right)^{-1} \cdot \tilde{\mathbf{T}}^{qp} \right) \cdot \mathbf{e}^{p*}$$

after the coarse grid correction we find

$$\mathbf{e}^p = \left(\mathbf{I}_p - \mathbf{T}^{pq} \cdot \left(\mathbf{L}^q \right)^{-1} \cdot \tilde{\mathbf{T}}^{qp} \cdot \mathbf{L}^p \right) \cdot \mathbf{e}_r^{p*}$$

Multigrid strategy

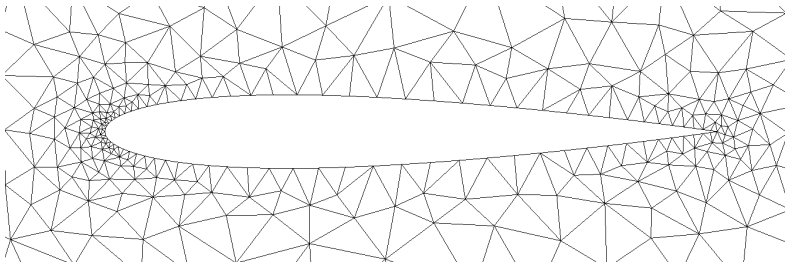
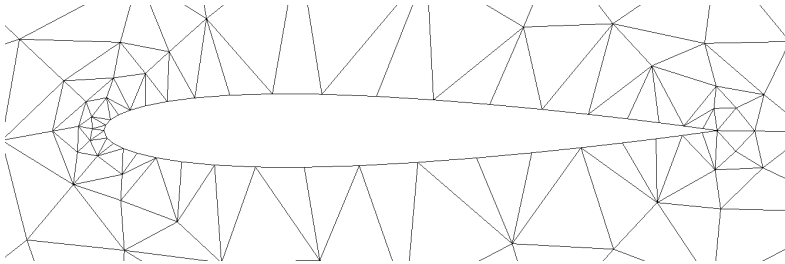
- cycling



- solvers and solver parameters for each level
 - explicit methods on fine levels
 - implicit methods on coarse levels

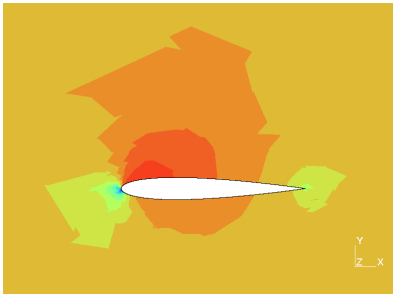
NACA0012 airfoil

$$M = 0.3, \alpha = 2^\circ$$

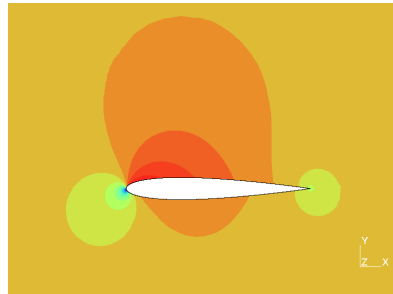


NACA0012: Mach isolines

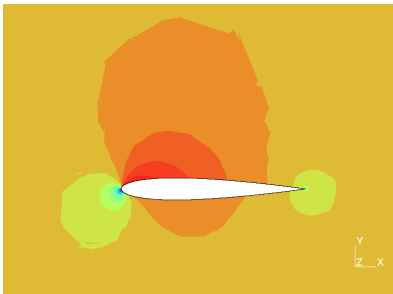
$p=1$, coarse



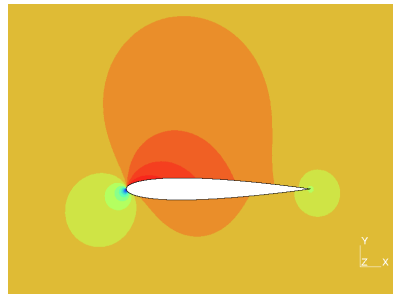
$p=4$, coarse



$p=1$, fine

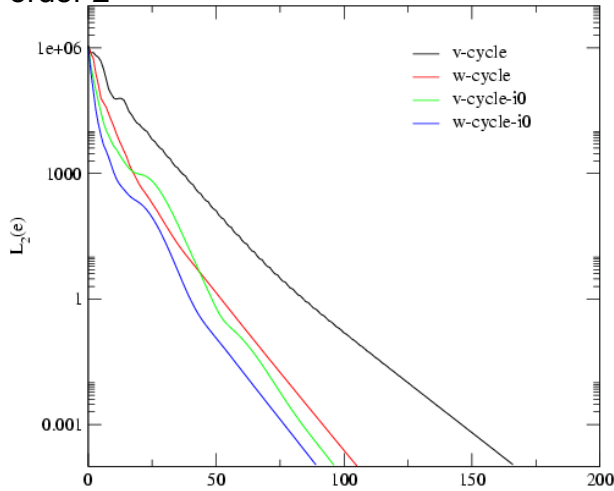


$p=4$, fine

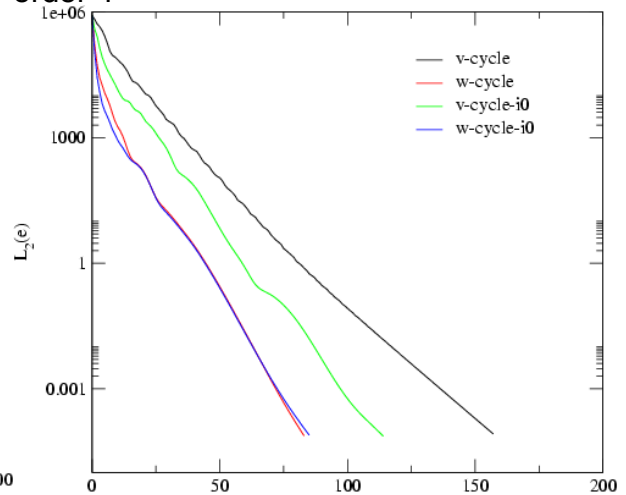


NACA0012 : strategies compared (cycles,coarse)

order 2

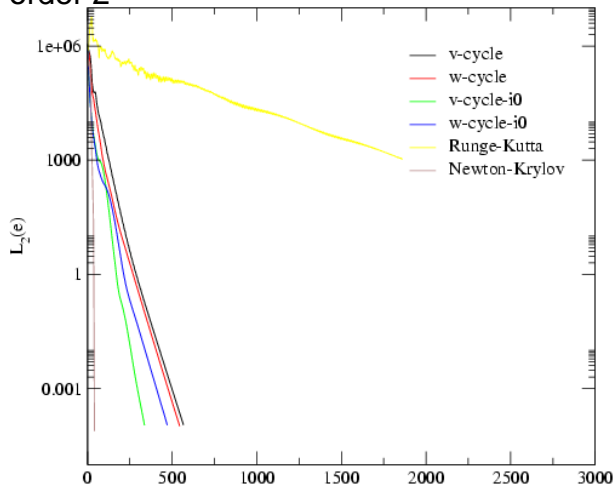


order 4

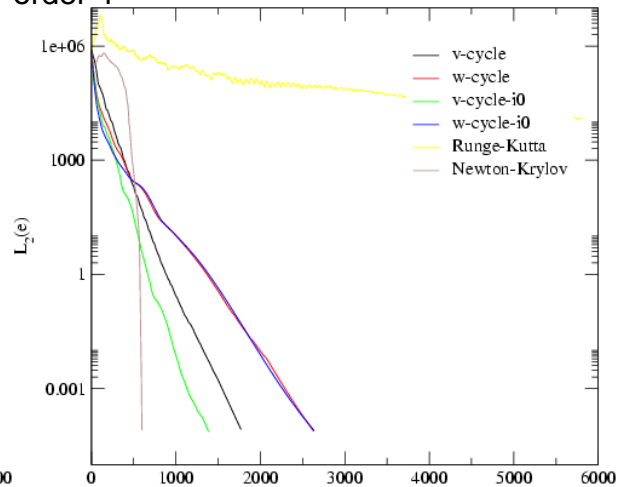


NACA0012 : strategies compared (CPU,coarse)

order 2

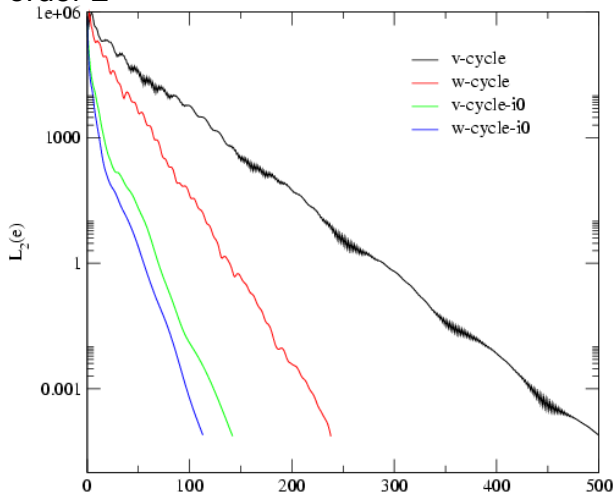


order 4

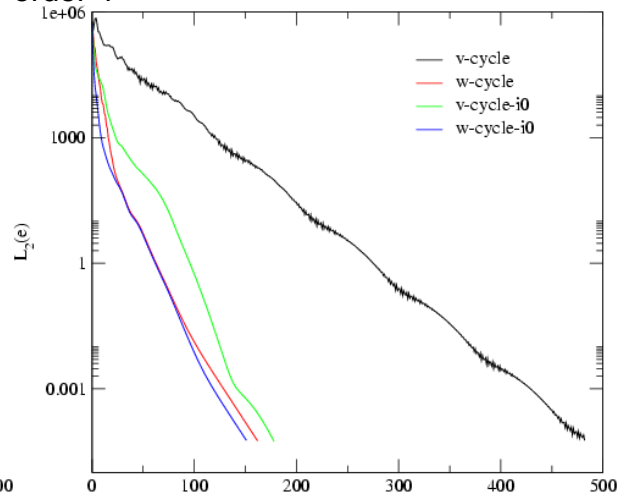


NACA0012 : strategies compared (cycles,fine)

order 2

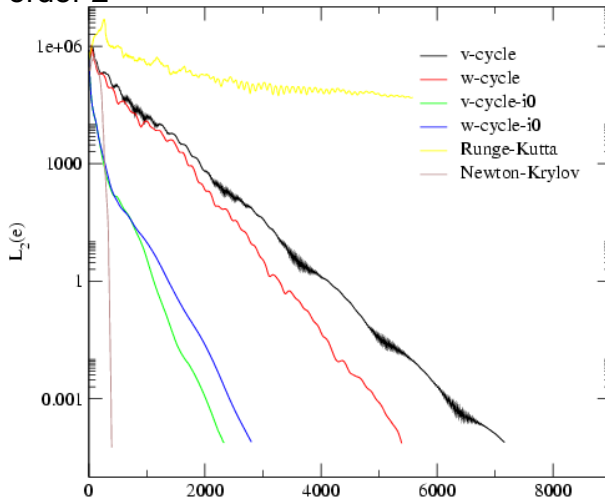


order 4

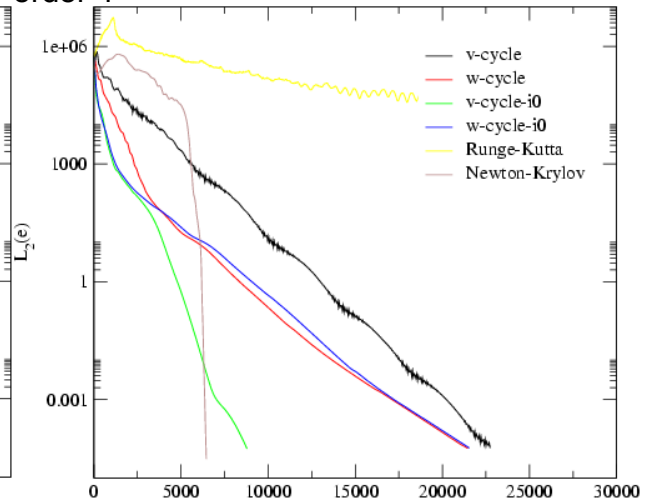


NACA0012 : strategies compared (CPU,fine)

order 2

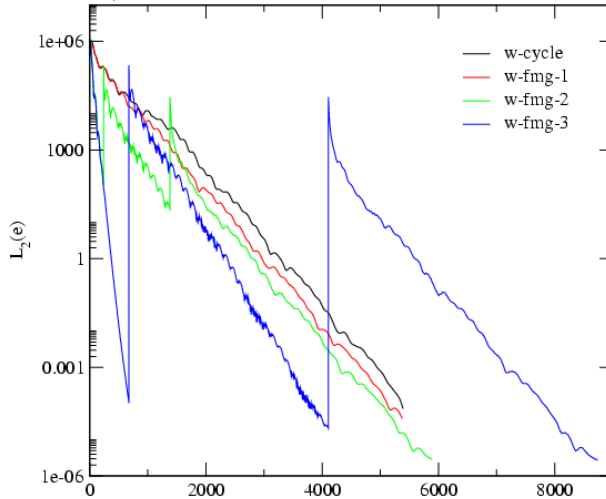


order 4

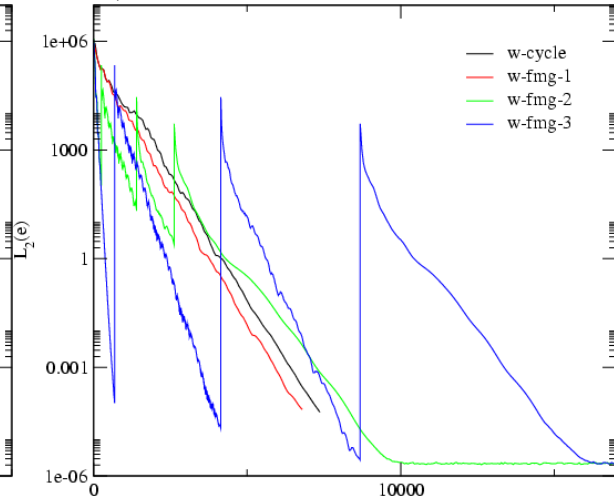


NACA0012 : strategies compared (FMG)

order 2, fine mesh



order 3, fine mesh



Conclusions and future work

Conclusions:

- generic framework for multigrid in DGFEM (equally applicable to h-multigrid on conforming or non-conforming meshes and non-hierarchic p-multigrid);
- p-MG convergence independent of order (TME), but dependent on grid size;
- using pure explicit smoothers, w-cycle is most performant;
- using implicit coarsest level v-cycle provides significant speed-ups, providing viable alternative to Newton-Krylov-ILU(0);
- FMG does not improve convergence rates.

Future work:

- look for a more performant smoother for higher order levels;
- combine with h-multigrid to obtain h- and p-independent convergence rates;
- use MG as preconditioner for GMRES;
- directional p-coarsening in boundary layers.

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